

A second order finite element method with mass lumping for Maxwell's equations on tetrahedra

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Maxwell's equations

Electromagnetic wave propagation in linear and non-dispersive but possibly inhomogeneous and anisotropic media

$$\begin{aligned}\varepsilon \partial_t E(t) &= \operatorname{curl} H(t) - \sigma E(t) && \text{in } \Omega \\ \mu \partial_t H(t) &= -\operatorname{curl} E(t) && \text{in } \Omega\end{aligned}$$

in Ω , with $E(0) = E_0$ and $H(0) = H_0$ in Ω and $n \times E(t) = 0$ on $\partial\Omega$

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Goal: systematic and flexible space discretization

- ▶ stable: no artificial energy production
- ▶ accurate: provable convergence rates
- ▶ efficient: appropriate for explicit time-stepping methods

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Methods: FDTD/FIT, FEM, FVM, DG, ...

Finite differences (FDTD/FIT)

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- ▶ 1966 - Yee - Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media
- ▶ 1977 - Weiland - Eine Methode zur Lösung der Maxwell'schen Gleichungen für sechskomponentige Felder auf diskreter Basis
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Pros

- ▶ Easy to implement
- ▶ stable, accurate $O(h^2 + \tau^2)$, efficient

Cons

- ▶ Difficulties in dealing with complex domains

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Approximation spaces: $V_h \subset H_0(\operatorname{curl}, \Omega)$ and $Q_h \subset L^2(\Omega)$

Galerkin method: For $t > 0$, find $E_h(t) \in V_h$ and $H_h(t) \in Q_h$ such that

$$\begin{aligned}(\varepsilon \partial_t E_h(t), v_h)_\Omega - (H_h(t), \operatorname{curl} v_h)_\Omega &= 0 \\ (\mu \partial_t H_h(t), q_h)_\Omega + (\operatorname{curl} E_h(t), q_h)_\Omega &= 0\end{aligned}$$

for all test functions $v_h \in V_h$ and $q_h \in Q_h$, and for all $t > 0$.

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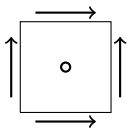
for all test functions $v_h \in V_h$ and $q_h \in Q_h$, and for all $t > 0$.

Algebraic realization. For a choice of basis functions, we have

$$\begin{aligned}\mathbf{M}_\varepsilon \partial_t \mathbf{e}(t) - \mathbf{C}^\top \mathbf{h}(t) &= 0 \\ \mathbf{D}_\mu \partial_t \mathbf{h}(t) + \mathbf{C} \mathbf{e}(t) &= 0\end{aligned}$$

Finite element spaces on reference elements.

► 1980 - Nedelec - Mixed Finite Elements in \mathbb{R}^3



$$V_h(Q) = \mathcal{N}_0(Q)$$

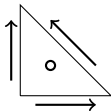
$$Q_h(Q) = P_0(Q)$$

$$\phi_1 = (1 - y, 0)$$

$$\phi_2 = (y, 0)$$

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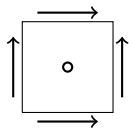
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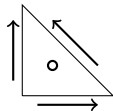
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Lemma (accuracy) If E and H are sufficiently smooth, then

$$\|E(t) - E_h(t)\|_{L^2} + \|H(t) - H_h(t)\|_{L^2} \leq Ch$$

- ▶ 1992 - Monk - Analysis of a finite element method for Maxwell's equations
- ▶ 1993 - Monk - An analysis of Nedelec's method for spatial discretization of Maxwell's equations

First order elements

Stability and accuracy.

Lowest order MFEM yields stable and accurate approximation in space.

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Numerical solution. Time integration of resulting ode system

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Note. Here \mathbf{D}_μ is diagonal, but \mathbf{M}_ε does not have a sparse inverse!
Thus, explicit time-stepping for standard MFEM is not efficient.

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Remedy – Mass-lumping: replace \mathbf{M}_ε by approximation \mathbf{M}_ε^L such that

- ▶ \mathbf{M}_ε^L corresponds to positive definite matrix (stability)
- ▶ \mathbf{M}_ε^L is good approximation for \mathbf{M}_ε (accuracy)
- ▶ $(\mathbf{M}_\varepsilon^L)^{-1}$ can be applied efficiently (efficiency)

construction of \mathbf{M}_ε^L usually via numerical quadrature.

Mass lumping literature

- ▶ 1975 - Fried, Malkus - Finite element mass matrix lumping by numerical integration with no convergence rate loss

- ▶ 1999 - Kong, Mulder, Veldhuizen - Higher-order triangular and tetrahedral finite elements with mass lumping for solving the wave equation
- ▶ 2000 - Becache, Joly, Tsogka - An analysis of new mixed finite elements for the approximation of wave propagation models
- ▶ 2001 - Mulder - Higher-order mass-lumped finite elements for the wave equation
- ▶ 2002 - Cohen - Higher-Order Numerical Methods for Transient Wave Equations

- ▶ 2018 - Geevers, Mulder, Vegt - New higher-order mass-lumped tetrahedral elements for wave propagation modelling

Mass-lumping in H^1

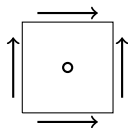
Mass lumping literature

- ▶ 1975 - Fried, Malkus - Finite element mass matrix lumping by numerical integration with no convergence rate loss
- ▶ 1990 - Lee, Madsen - A mixed FEM formulation for Maxwell's equations in the time domain
- ▶ 1995 - Cohen, Monk - Mass lumped edge elements in three dimensions
- ▶ 1997 - Elmkins, Joly - Elements finis d'arête et condensation de masse pour les équations de Maxwell - le cas 3D
- ▶ 1998 - Cohen, Monk - Gauss Point Mass Lumping Schemes for Maxwell's Equations
- ▶ 1999 - Kong, Mulder, Veldhuizen - Higher-order triangular and tetrahedral finite elements with mass lumping for solving the wave equation
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- ▶ 2002 - Cohen - Higher-Order Numerical Methods for Transient Wave Equations
- ▶ 2004 - Lacoste - Mass-lumping for the first order Raviart–Thomas–Nedelec finite elements
- ▶ 2007 - Jund, Salmon - Arbitrary high-order finite element schemes and high order mass lumping
- ▶ 2018 - Geevers, Mulder, Vegt - New higher-order mass-lumped tetrahedral elements for wave propagation modelling

Mass-lumping in H^1

Mass-lumping in $H(\text{div})$ and $H(\text{curl})$

Observation for the lowest order case



$$V_h(Q) = \mathcal{N}_0(Q)$$

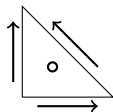
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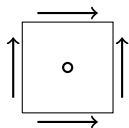
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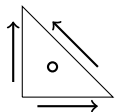
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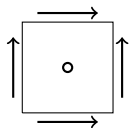
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Observation: No combination of **quadrature rule** and **basis functions** that leads to decoupling of entries in mass matrix for V_h .

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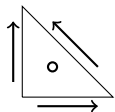
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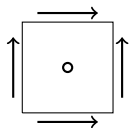
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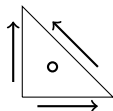
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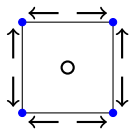
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Some existing methods: Acute mesh lumping (triangles)

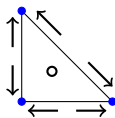
- ▶ 1996 - Baranger - Connection between finite volume and mixed finite element methods

First order elements... with mass lumping!

Use a larger polynomial space



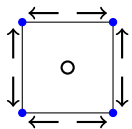
$$\begin{aligned}\tilde{V}_h(Q) &= \mathcal{NC}_1(Q) \\ \tilde{Q}_h(Q) &= P_0(Q)\end{aligned}$$



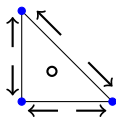
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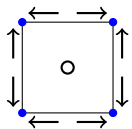


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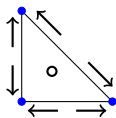
Lemma. $\tilde{\mathbf{M}}_\epsilon^L$ is block diagonal and thus also $(\tilde{\mathbf{M}}_\epsilon^L)^{-1}$.

First order elements... with mass lumping!

Use a larger polynomial space

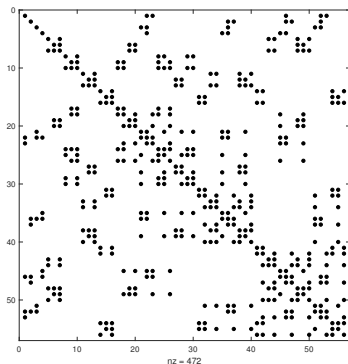
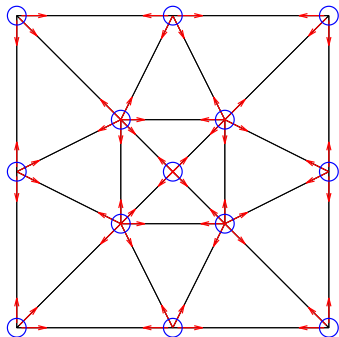


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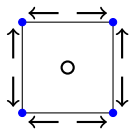
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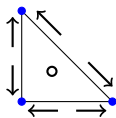


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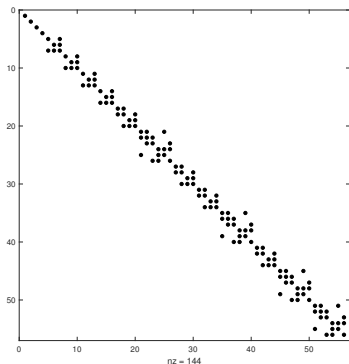
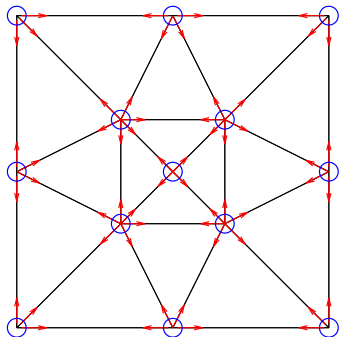


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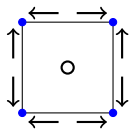
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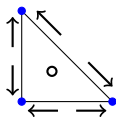


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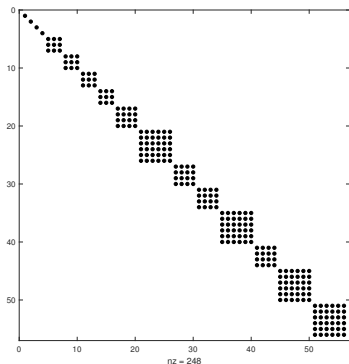
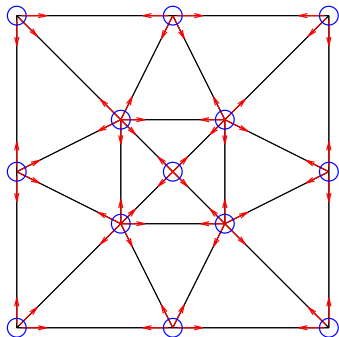


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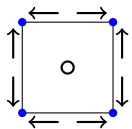
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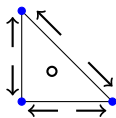


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Theorem (accuracy)

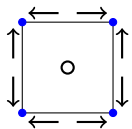
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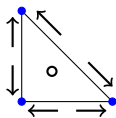
- ▶ 2020 - Egger, Radu - A mass-lumped mixed finite element method for Maxwell's equations

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$$\begin{aligned}\tilde{V}_h(T) &= \mathcal{NC}_1(T) \\ \tilde{Q}_h(T) &= P_0(T)\end{aligned}$$

Lemma. $\tilde{\mathbf{M}}_\epsilon^L$ is block diagonal and thus also $(\tilde{\mathbf{M}}_\epsilon^L)^{-1}$.

Theorem (accuracy)

If \mathbf{E} and \mathbf{H} are sufficiently smooth, then

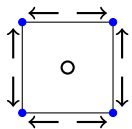
$$\|\mathbf{E}(t) - \tilde{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \tilde{\mathbf{H}}_h(t)\| \leq Ch$$

- ▶ 2020 - Egger, Radu - A mass-lumped mixed finite element method for Maxwell's equations

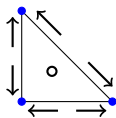
Proof Idea: Error splitting in discrete and projection error, discrete stability, energy estimates, consistency error, **analysis of the quadrature error (Strang)**.

First order elements... with mass lumping!

Use a larger polynomial space



$$\begin{aligned}\tilde{V}_h(Q) &= \mathcal{NC}_1(Q) \\ \tilde{Q}_h(Q) &= P_0(Q)\end{aligned}$$



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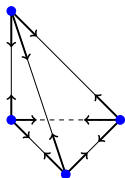
- ▶ 2020 - Egger, Radu - A mass-lumped mixed finite element method for Maxwell's equations

Proof Idea: Error splitting in discrete and projection error, discrete stability, energy estimates, consistency error, **analysis of the quadrature error (Strang)**.

Requirement : The quadrature rule must be exact for $P_0(T)^2 \times \tilde{V}_h(T)$

First order elements on tetrahedral meshes

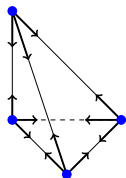
The same concept also applies in 3D on tetrahedral meshes



$$\begin{aligned}\tilde{V}_h(T) &= \mathcal{NC}_1(T) \\ \tilde{Q}_h(T) &= P_0(T)\end{aligned}$$

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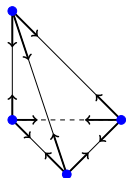
Theorem (accuracy)

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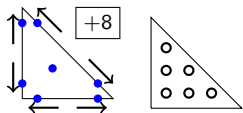
$$\|\mathbf{E}(t) - \tilde{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \tilde{\mathbf{H}}_h(t)\| \leq Ch$$

Next task : Extension to second order elements.

Second order elements

Extension to second order elements

- ▶ 1997 - Elmkies, Joly - Elements finis d'arete et condensation de masse pour les equations de Maxwell - le cas 2D

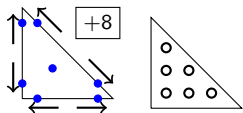


$$\widehat{V}_h(T) = \mathcal{N}_1(T) \oplus B = \mathcal{E}\mathcal{J}_1(T) \subseteq P_3(T)$$
$$\widehat{Q}_h(T) = P_2(T)$$

Second order elements

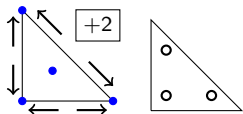
Extension to second order elements

- ▶ 1997 - Elmkies, Joly - Elements finis d'arete et condensation de masse pour les equations de Maxwell - le cas 2D



$$\begin{aligned}\widehat{V}_h(T) &= \mathcal{N}_1(T) \oplus B = \mathcal{E}\mathcal{J}_1(T) \subseteq P_3(T) \\ \widehat{Q}_h(T) &= P_2(T)\end{aligned}$$

New proposal :

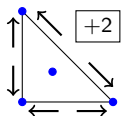


$$\begin{aligned}\widehat{V}_h(T) &= \mathcal{N}_1(T) \subseteq P_2(T) \\ \widehat{Q}_h(T) &= P_1(T)\end{aligned}$$

The quadrature rule is exact for P_2 polynomials ... but is this enough ?

Short notes on the analysis

New proposal :



$$\widehat{V}_h(T) = \mathcal{N}_1(T)$$

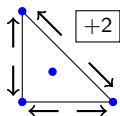
$$\widehat{Q}_h(T) = P_1(T)$$

Theorem (accuracy). If \mathbf{E} and \mathbf{H} are sufficiently smooth, then

$$\|\mathbf{E}(t) - \widehat{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \widehat{\mathbf{H}}_h(t)\| \leq Ch^2$$

Short notes on the analysis

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$$\widehat{V}_h(T) = \mathcal{N}_1(T)$$

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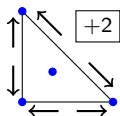
$$\|\mathbf{E}(t) - \widehat{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \widehat{\mathbf{H}}_h(t)\| \leq Ch^2$$

Proof Idea: Discrete stability, energy estimates, Galerkin orthogonality, consistency error, Strang analysis of the quadrature error.

Classic requirement : The quadrature rule has to be exact for $P_1(T)^d \times \widehat{V}_h(T)$

Short notes on the analysis

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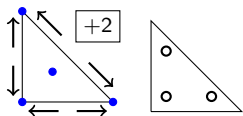
Classic requirement : The quadrature rule has to be exact for $P_1(T)^d \times \widehat{V}_h(T)$

New requirements

- (i) There exists a splitting $\widehat{V}_h(T) = W(T) \oplus B(T)$ such that $\dim(B(T)) = \dim(\text{curl}(B(T)))$ and $\text{curl}(B(T)) \cap \text{curl}(W(T)) = \{0\}$
- (ii) The quadrature rule is exact for $P_1(T)^2 \times W(T)$

Short notes on the analysis

New proposal :



$$\begin{aligned}\widehat{V}_h(T) &= \mathcal{N}_1(T) = \mathcal{NC}_1(T) \oplus B(T) \\ \widehat{Q}_h(T) &= P_1(T)\end{aligned}$$

Theorem (accuracy). If \mathbf{E} and \mathbf{H} are sufficiently smooth, then

$$\|\mathbf{E}(t) - \widehat{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \widehat{\mathbf{H}}_h(t)\| \leq Ch^2$$

Proof Idea: Discrete stability, energy estimates, Galerkin orthogonality, consistency error, Strang analysis of the quadrature error.

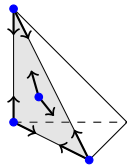
Classic requirement : The quadrature rule has to be exact for $P_1(T)^d \times \widehat{V}_h(T)$

New requirements

- (i) There exists a splitting $\widehat{V}_h(T) = \mathcal{NC}_1(T) \oplus B(T)$ such that $\dim(B(T)) = \dim(\text{curl}(B(T)))$ and $\text{curl}(B(T)) \cap \text{curl}(\mathcal{NC}_1(T)) = \{0\}$
- (ii) The quadrature rule is exact for $P_1(T)^2 \times \mathcal{NC}_1(T) = P_2(T)^2$

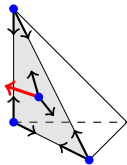
Second order method - 3D

- ▶ 1997 - Elmkies, Joly - Elements finis d'arete et condensation de masse pour les equations de Maxwell - le cas 3D



$$\widehat{V}_h(T) = \mathcal{N}_1(T)$$

- ▶ 1997 - Elmkies, Joly - Elements finis d'arete et condensation de masse pour les equations de Maxwell - le cas 3D

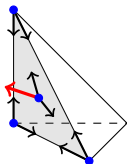


$$\widehat{V}_h(T) = \mathcal{N}_1(T) \oplus B(T)$$

$$B(T) = \text{span} \left\{ \begin{array}{l} \Phi_1 = \lambda_2 \lambda_3 \lambda_4 \nabla \lambda_1 \\ \Phi_2 = \lambda_1 \lambda_3 \lambda_4 \nabla \lambda_2 \\ \Phi_3 = \lambda_1 \lambda_2 \lambda_4 \nabla \lambda_3 \\ \Phi_4 = \lambda_1 \lambda_2 \lambda_3 \nabla \lambda_4 \end{array} \right.$$

Second order method - 3D

- ▶ 1997 - Elmkies, Joly - Elements finis d'arete et condensation de masse pour les equations de Maxwell - le cas 3D



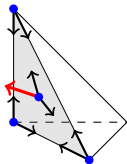
$$\widehat{V}_h(T) = \mathcal{N}_1(T) \oplus B(T) \quad B(T) = \text{span} \left\{ \begin{array}{l} \Phi_1 = \lambda_2 \lambda_3 \lambda_4 \nabla \lambda_1 \\ \Phi_2 = \lambda_1 \lambda_3 \lambda_4 \nabla \lambda_2 \\ \Phi_3 = \lambda_1 \lambda_2 \lambda_4 \nabla \lambda_3 \\ \Phi_4 = \lambda_1 \lambda_2 \lambda_3 \nabla \lambda_4 \end{array} \right.$$

But $\nabla(\lambda_1 \lambda_2 \lambda_3 \lambda_4) = \Phi_1 + \Phi_2 + \Phi_3 + \Phi_4 \Rightarrow \text{curl}(\Phi_1 + \Phi_2 + \Phi_3 + \Phi_4) = 0$.

Thus $\dim(B(T)) \neq \dim(\text{curl } B(T))$!

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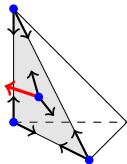
Theorem [EggerRadu21]. If (and only if) $\text{div}(\mathbf{E}) = 0$, then

$$\|\mathbf{E}(t) - \widehat{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \widehat{\mathbf{H}}_h(t)\| \leq Ch^2$$

Note. In general, second order convergence is lost!

Second order method - 3D

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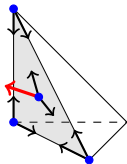
$$\|\mathbf{E}(t) - \widehat{\mathbf{E}}_h(t)\| + \|\mathbf{H}(t) - \widehat{\mathbf{H}}_h(t)\| \leq Ch^2$$

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Solution. Modify one basis function, for example $\widehat{\Phi}_4 = \lambda_1 \lambda_2 \lambda_3 (\lambda_2 - \lambda_1 + 1) \nabla \lambda_4$

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Main takeaways

Key ingredients for mass lumping:

- ▶ Start with a basis space V_h that contains all $P_k(T)^d$ polynomials (for approximation).
- ▶ V_h dictates the number of continuity conditions on the boundary
- ▶ Find a quadrature rule that has sufficiently many quadrature points on the boundary and has the desired accuracy
- ▶ Extend V_h by appropriate "bubble" functions such that we have exactly d -many functions for each quadrature point.

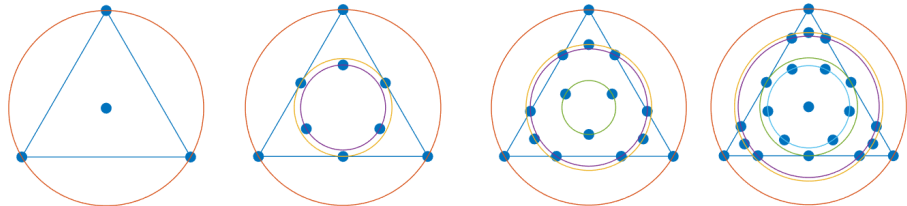
List of relevant publications

- ▶ 2020 - Egger, Radu - A mass-lumped mixed finite element method for acoustic wave propagation.
- ▶ 2020 - Egger, Radu - A mass-lumped mixed finite element method for Maxwell's equations
- ▶ 2021 - Egger, Radu - A second order finite element method with mass lumping for wave equations in $H(\text{div})$.
- ▶ 2021 - Egger, Radu - A Second-Order Finite Element Method with Mass Lumping for Maxwell's Equations on Tetrahedra.

Thank you for your attention!

Extension to even higher orders

We look for Gauss-Lobatto type quadrature rules !



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